Patterns in Geomorphology

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Patterns of erosion and chemical dissolution due to running water from snow melting (Lapiaz)



Patterns of erosion and chemical dissolution



Erosion patterns in badlands (Tzin valley, Israel)

A BARRAN A STATEMENT

Patterns of erosion and deposition: Meandering rivers (Ucayali river, Peru)





Patterns of erosion and deposition: Braided rivers

River networks (Yemen)

Coastal patterns (coast of Carolina, USA)



Courtesy of A. Brad Murray



Aeolian patterns Great Sand Dunes National Monument, Colorado, photo by Bob Bauer

Aeolian bedforms in deserts and sandy beaches: Ripples (length of cm, amplitude of mm) Megaripples (length of meters, amplitude of cm) Dunes (length of tens or hundreds of meters, amplitude of (tens of) meters)



Southern Negev desert, Israel, photo by Hezi Yizhaq

Types of (non-vegetated) dunes:

Barchans

Transverse

Seif (linear) dunes

Star dunes

Barchan dunes on Mars



Namib desert

Linear dunes, Mauritania



Namib sand sea



Complex dune structures (Saudi Arabia)



Complex dune structures (Gran Desierto))



Superposed aeolian bedforms



Blown by wind: Nonlinear dynamics of aeolian sand ripples

Properties of aeolian ripples: Ripple index about 15-20 Almost 1D bedforms: crests are perpendicular to the wind, with defects Slight asymmetry between lee and stoss slopes (no slip face) Need of a wind intensity threshold to have ripple formation Rapid response to wind The ripple pattern coarsens with time and it slowly moves downstream Grain sorting



A short history of ripple studies:

R.A. Bagnold, *The Physics of Blown Sand and Desert Dunes*, 1941
R. Cooke, A. Warren, A. Goudie, *Desert Geomorphology*, 1993
N. Lancaster, *Geomorphology of desert dunes*, 1995

R.P. Sharp, J. of Geology, 1963
M. Seppala and K. Lindé, Geografiska Annaler, 1978
B.B. Willetts and M.A. Rice, 1983-1989
J.E. Ungar and P.K. Haff, Sedimentology, 1987
R.S. Anderson, Sedimentology, 1987
R.S. Anderson, Earth Sci. Rev., 1990

B.T. Werner and D.T. Gillespie, *PRL*, 1993W. Landry and B.T. Werner, *Physica D*, 1994H. Nishimori and N. Ouchi, *PRL*, 1993

R.B. Hoyle and A.W. Woods, *PRE*, 1997 L. Prigozhin, *PRE*, 1999

O. Terzidis, P. Claudin and J.-P. Bouchaud, *Eur. Phys. J. B*, 1998
A. Valance and F. Rioual, *Eur. Phys. J. B*, 1999
Z. Csahok and C. Misbah, *Eur. Phys. J. E*, 2000
Z. Csahok, C. Misbah, F. Rioual and A. Valance, *cond-mat*, 2000
H. Yizhaq, N.J. Balmforth, A. Provenzale, Physica D, 2004
H. Yizhaq et al. EPS1, 2019 (megaripples)

Mechanism of ripple formation: aeolian ripples form due to the instability of a flat sand bed exposed to strong wind

When the wind starts to blow, sand grains are lifted into the air. These grains are accelerated by the wind and fall down, hit the surface, and eject other grains.

The rebounding (saltating) grains are then accelerated by the wind, and a cascade process ensues.

An entire population of saltating grains emerges. The height of the saltation layer can be about one meter in strong winds Hypotheses of ripple formation (Cooke et al, 1993):

1. A rithmic barrage of saltating grains (Bagnold)

2. The wave hypothesis:
a) The bed as a fluid
b) The saltation curtain as a fluid
c) Wave-like instabilities in the boundary layer
d) Secondary motions in the lee of transverse ripples

3. The role of reptating grains (Anderson 1987)



Saltation with typical jump length *L* Reptation with typical jump length *a* << *L L* is about *1 m* and *a* is about *1 cm*



Upon impact, the energy of a saltating grain goes as: 80% to one (on average) rebounding grain, 10% to a few reptating grains, 10% to the deformation of the bed



Depth of the saltation layer: up to about 1 m Depth of the reptation layer: a few mm We can idealize the problem in terms of a sand surface bombarded by a continuous flux of saltating grains that hit the surface at constant (small) angle $\phi = 8-12^{\circ}$

> The saltating grains drive the system. The important dynamics is contained in the behavior of the reptating grains

Conservation of sand

$$(1-n)\rho\frac{\partial\zeta}{\partial t} = -\nabla\cdot\mathbf{Q}$$

- *n* porosity of the bed (about 0.35)
- ρ density of sand
- ζ elevation of the sand surface
- **Q** flux of sand grains

The flux of sand:

$\mathbf{Q} = \mathbf{Q}_s + \mathbf{Q}_r$

- Q_s flux of saltating grains
 Q flux of rontating grains
- \mathbf{Q}_r flux of reptating grains

The flux of saltating grains is assumed to be constant for aeolian ripples

 $\nabla \cdot \mathbf{Q}_s = 0$

All the dynamics is contained in the variability of the reptation flux

NB: This is untenable for megaripples and dunes

There is no feedback of the aeolian bedforms on the wind and on the flux of saltating grains

NB: This is untenable for dunes

1D case

if all the grains had the same reptation length a

$$Q_r^{bare}(x,t) = m \int_{x-a}^{x} N_{ej}(x') dx'$$

- *m* mass of a sand grain
- N_{ej} number of ejected grains
- N_r average number of reptating grains ejected by one saltating grain
- $N_{ej} = N_r N_{im}$ where N_{im} is the number of impacting grains

For a distribution of reptation lengths

$$Q_r^{bare}(x,t) = m \int d\alpha \ p(\alpha) \int N_{ej}(x') dx'$$

 $x - \alpha$

- *m* mass of a sand grain
- N_{ei} number of ejected grains

 $-\infty$

- N_r average number of reptating grains ejected by one saltating grain
- $N_{ej} = N_r N_{im}$ where N_{im} is the number of impacting grains
- $p(\alpha)$ distribution of reptation lengths $\int \alpha p(\alpha) d\alpha = a$

Important angles



Number density of impacting grains

$$N_{im}(x) = N_{im}^0 \left(1 + \frac{\tan \theta}{\tan \phi} \right) \cos \theta$$

 $= N_{im}^0 \cot \phi \ \frac{\tan \phi + \zeta_x}{\sqrt{1 + \zeta_x^2}}$

 N_{im}^{0} number density of impacting grainson a horizontal surface (about $10^7 \text{ m}^{-2} \text{ s}^{-1}$) θ inclination of the surface

 $tan \theta = \zeta_x$

$$a = 2\frac{V^2}{g}\frac{\sin(\gamma - \theta)\cos\gamma}{\cos\theta} = a_{hor}(1 - \cot\gamma\tan\theta)$$



Shadowing: the flux of reptating grains becomes

$N_{im}(x) = N_{im}^{0} \cot \phi Max \left\{ \frac{\tan \phi + \zeta_x}{\sqrt{1 + \zeta_x^2}}, 0 \right\}$

The full integral model:

 $\frac{\partial \zeta}{\partial t} = -Q_0 \frac{\partial}{\partial x} \left[(1 - \mu \zeta_x) \int_{-\infty}^{\infty} d\alpha \, p(\alpha) \int_{x-\alpha}^{x} F(x') \, dx' \right]$

$$F(x) = Max \left\{ \frac{\tan \phi + \zeta_x}{\sqrt{1 + \zeta_x^2}}, 0 \right\}$$

$$Q_0 = \frac{mN_r N_{im}^0 \cot \phi}{\rho(1-n)}$$

Linear stability analysis $Q_r(x,t) = Q_r^{bare}(x,t) (1 - \mu \zeta_x)$





height (m)

Time Evolution of Ripples

Coarsening of the ripple pattern





An extension to the 2D case:

If all the grains had the same jump length a, the flux at (x, y)in the direction ψ to the x axis is assumed to be proportional to the number of grains ejected between $(x - a \cos \psi, y - a \sin \psi)$ and (x, y)



Fig. 12. Snapshots of a two-dimensional ripple field produced by the numerical integration of Eq. (28) on a 256 × 256 grid with periodic boundary condition, at T = 2, 4, 8, 16. The wind direction is from left to right. Parameter values: $\epsilon = \mu = 0.25, \phi = 10^{\circ}, \beta_1 = \beta_3 = 1, \beta_2 = \tan \phi$ and a domain size of $5\pi/2$.





Megaripples on Earth and Mars

The origin of the transverse instability of aeolian megaripples

H. Yizhaq^{a,*}, G. Bel^{a,b}, S. Silvestro^{c,d}, T. Elperin^e. I.F. Kok^f. M. Cardinale^g. A. Provenzale^h. I. Katraⁱ Earth and Planetary Science Letters 512 (2019) 59–70





Field work



Laboratory experiments



Numerical simulations



Numerical simulations



Numerical simulations



Lateral flux dependent on height





Mars case

A much harder problem: The dynamics of aeolian sand dunes Emma Pike - Originally uploaded to English Wikipedia



Just a beginning in the fascinating world of (eco)geomorphological pattern modelling

26 October 2007 (UTC) - Own worl ex.php?curid=2978390